

FINITE ELEMENTS IN FLUIDS

ASSIGNMENT2- 1D UNSTEADY TRANSPORT PROBLEM

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LEAP FROG METHOD

The galerkin formulation for the leap-frog method is presented below-

$$\frac{u^{n+1} - u^{n-1}}{2\Delta t} = u_t^n, \quad u_t = -a \cdot \nabla u$$

$$\Rightarrow \frac{u^{n+1} - u^{n-1}}{2\Delta t} = (-a \cdot \nabla u^n)$$

$$\Rightarrow u^{n+1} - u^{n-1} = -2a\Delta t \nabla u^n$$

\Rightarrow Using galerkin formulation, we have

$$(w, u^{n+1}) - (w, u^{n-1}) = -2a\Delta t (w, \nabla u^n)$$

$$\Rightarrow [w, (u^{n+1} - u^{n-1})] = -2a\Delta t (w, \nabla u^n)$$

Changes in the code (System.m) made to accommodate this code is given below-

```
21 - case 5 % Leap-Frog
22 -     A = M;
23 -     B = -2*a*dt*C;
24 -     methodName = 'LF';
```

The following changes were made to the main.m

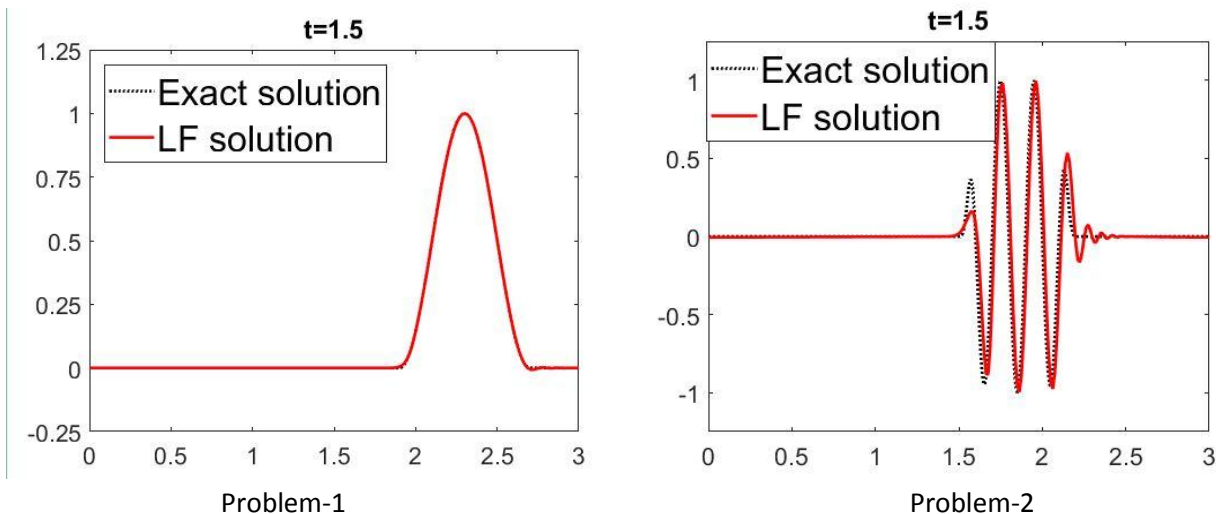
```

86 -     if method==5
87 -     □   for n= 1:nStep
88 -         if n==1
89 -             [A,B,methodName]= System(1,M,K,C,a,dt);
90 -             Du = A \ (B*u(1:nPt,n) + f);
91 -             u(1:nPt,n+1) = u(1:nPt,n) + Du;
92 -             clear A,B;
93 -         else
94 -             [A,B,methodName]= System(5,M,K,C,a,dt);
95 -             Du = A \ (B*u(1:nPt,n) + f);
96 -             u(1:nPt,n+1) = u(1:nPt,n-1) + Du;
97 -         end
98 -     end

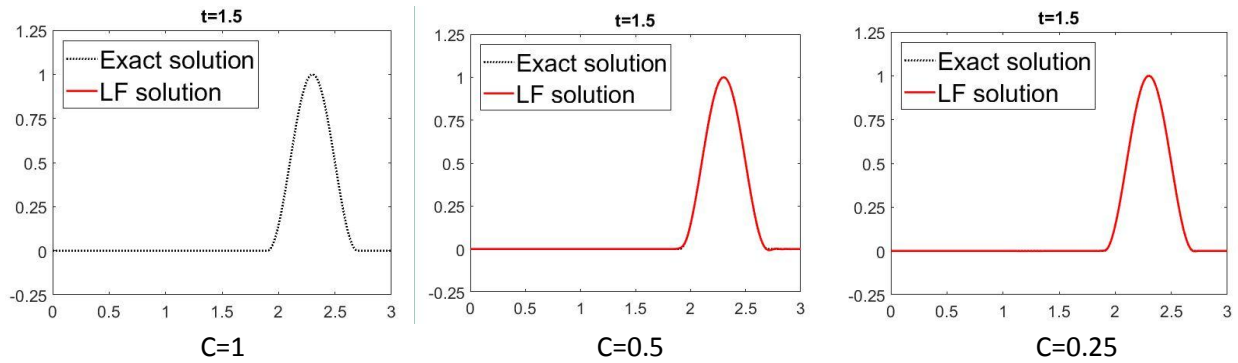
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It should be noted that as for $n=1$, $u(n-1)$ doesn't exist, Leap frog formulation is ineffective. So we need to implement another method to initialize the Leap-Frog then from $n=2$, Leap frog takes over further computation. Here I have used Lax-Wendroff + Galerkin formulation to initialize the Leap-Frog formulation.

Results obtained-



The simulation has been done for 2 problems $C=0.5$ and the code seems to be working. To check the stability of the method the following results were obtained by varying the courant number for problem 1 which seems to be consistent with the theory. Leap frog method is stable for courant number $C^2 \leq (3/4)$, which is consistent with the results. It is unstable for $C=1$ but stable for values less than 0.8.



THIRD ORDER TAYLOR-GALERKIN

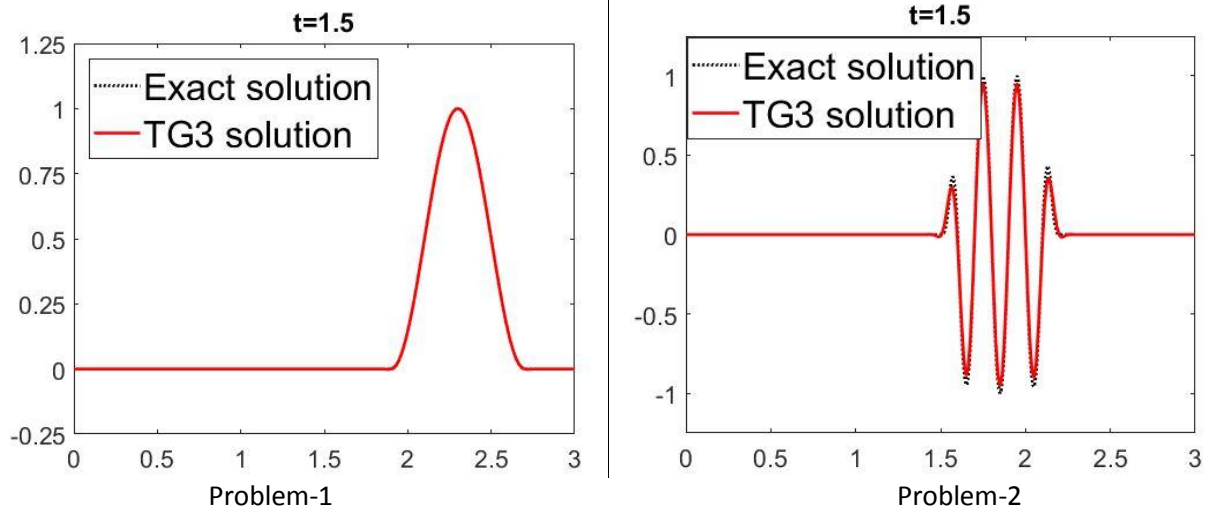
As per the theory given in the slides, the following changes were made to the system.m file-

```

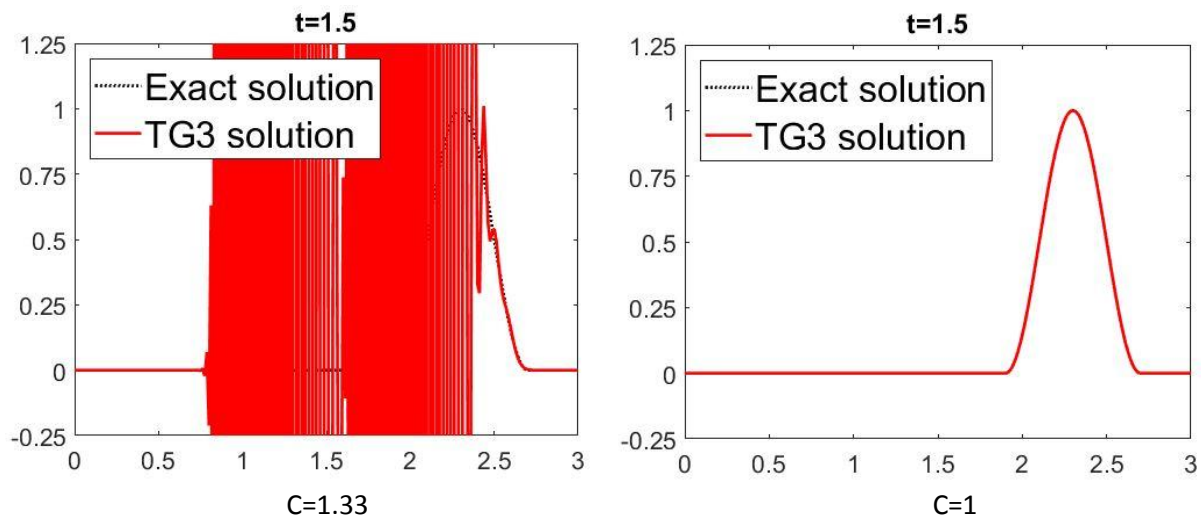
25 - case 6 % TG3
26 -     A = M + 0.16667*dt^2*a^2*K;
27 -     B = -a*dt*C- 0.5*a^2*dt^2*K;
28 -     methodName = 'TG3';

```

Results obtained-



The simulation has been done for 2 problems C=0.5 and the code seems to be working. To check the stability of the method the following results were obtained by varying the courant number for problem 1 which seems to be consistent with the theory. TG3 method is stable for courant number $C^2 \leq 1$, which is consistent with the results. It is unstable for $C > 1$ but stable for values less than equal to 1.



THIRD ORDER TAYLOR-GALERKIN 2-STEP

The galerkin formulation for the TG3-2S method is presented below-

First step

$$\tilde{u}^n = u^n + \frac{1}{3} \Delta t u_t^n + \alpha \Delta t^2 u_{tt}^n$$

$$\Rightarrow \tilde{u}^n = u^n + \frac{1}{3} \Delta t u_t^n + \frac{1}{9} \Delta t^2 u_{tt}^n \quad [\alpha = 1/9]$$

$$\Rightarrow \tilde{u}^n = u^n + \frac{1}{3} \Delta t (-a \cdot \nabla u^n) + \frac{1}{9} \Delta t^2 (-a \cdot \nabla)^2 u^n$$

$$\Rightarrow \tilde{u}^n = u^n + \left(-\frac{a \Delta t}{3}\right) \nabla u^n + \frac{\Delta t^2}{9} a^2 \nabla^2 u^n$$

Using Galerkin formulation,

$$(\omega, \tilde{u}^n) = (\omega, u^n) - \frac{\Delta t}{3} a (\omega, \nabla u^n) + \frac{a^2 \Delta t^2}{9} (\omega, \nabla^2 u^n)$$

$$[\omega, \nabla^2 u^n]_{\Omega} = [\omega, \nabla u^n]_{\Gamma}^0 - [\nabla \omega, \nabla u^n]_{\Omega} \quad [\text{Boundary flux in this case is neglected}]$$

$$\therefore (\omega, \tilde{u}^n) = (\omega, u^n) - \frac{a \Delta t}{3} (\omega, \nabla u^n) - \frac{a^2 \Delta t^2}{9} (\nabla \omega, \nabla u^n)$$

From here we can get the value for \tilde{u}^n and substitute it in the second step of formulation which is presented below-

This can be clearly seen from the two step code in (main.m) given below-

Second step

$$u^{n+1} = u^n + \Delta t \tilde{u}_t^n + \frac{1}{2} \Delta t^2 \tilde{u}_{tt}^n$$

$$\Rightarrow u^{n+1} = u^n + \Delta t (-a \cdot \nabla u^n) + \frac{1}{2} \Delta t^2 (-a \cdot \nabla)^2 \tilde{u}^n$$

$$\Rightarrow u^{n+1} = u^n - a \Delta t \nabla u^n + \frac{\Delta t^2 a^2}{2} \nabla^2 \tilde{u}^n$$

Using galerkin formulation,

$$(\omega, u^{n+1}) = (\omega, u^n) - a \Delta t (\omega, \nabla u^n) + \frac{a^2 \Delta t^2}{2} (\omega, \nabla^2 \tilde{u}^n)$$

$$[\omega, \nabla^2 \tilde{u}^n]_{\Omega} = [\omega, \nabla \tilde{u}^n]_{\Gamma} - [\nabla \omega, \nabla \tilde{u}^n]$$

$$\Rightarrow \underbrace{[\omega, (u^{n+1} - u^n)]}_{\Delta u} = -a \Delta t (\omega, \nabla u^n) + \frac{a^2 \Delta t^2}{2} (-\nabla \omega, \nabla \tilde{u}^n)$$

Changes in the code (System.m) made to accommodate this code is given below- TG3-I and TG3-II represent the two steps-

```

29 -     case 7 % 2 step TG3-I
30 -         A = M;
31 -         B = -(1/3)*dt*a*C-(1/9)*dt^2*a^2*K;
32 -         methodName = 'TG3';
33 -     case 8 % 2 step TG3-II
34 -         A = M;
35 -         B = -a*dt*C;
36 -         methodName = 'TG3';

```

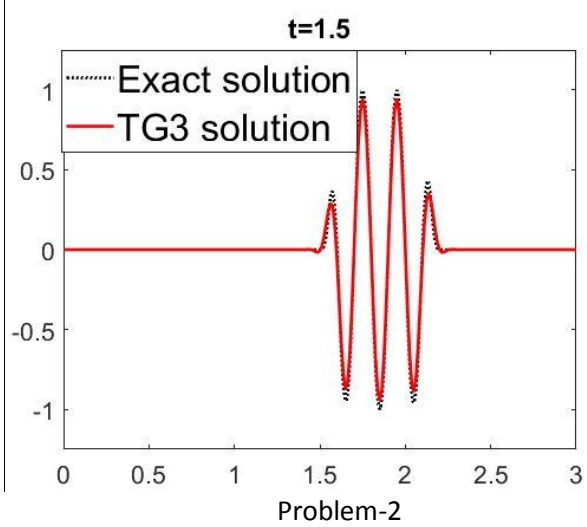
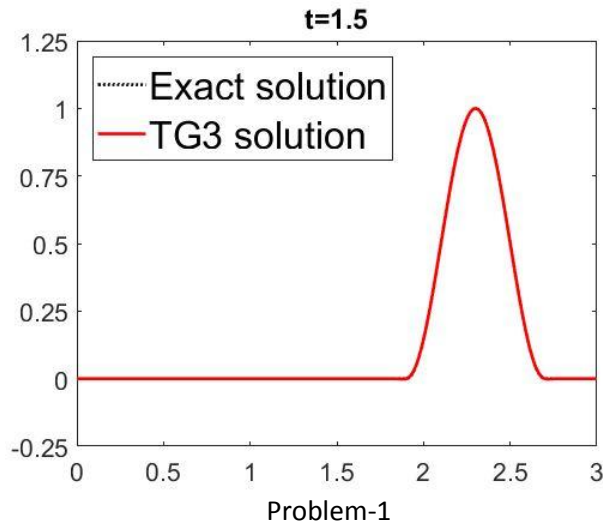
The following changes were made to the main.m

```

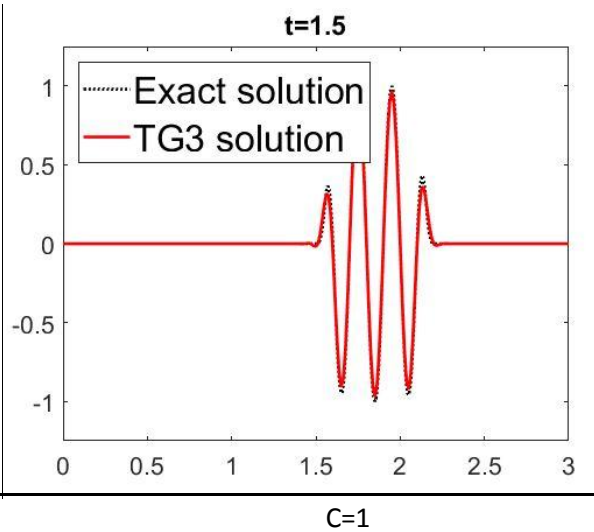
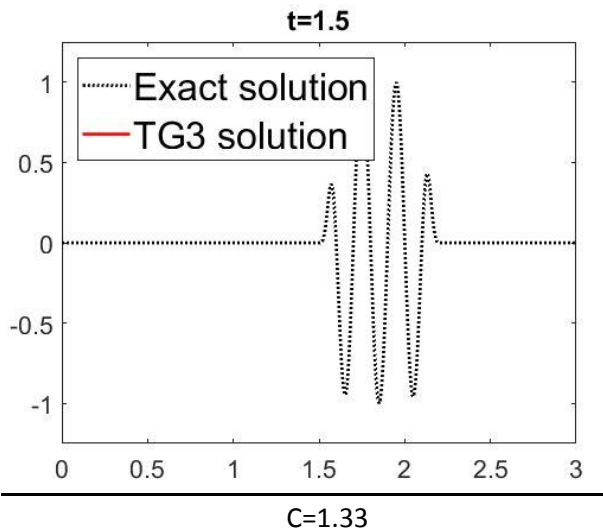
99 -     else if method==7
100 -         for n= 1:nStep
101 -             [A,B,methodName]= System(7,M,K,C,a,dt);
102 -             Du = A\ (B*u(1:nPt,n) + f);
103 -             u_bar= u(1:nPt,n) + Du;
104 -             clear A,B;
105 -             [A,B,methodName]= System(8,M,K,C,a,dt);
106 -             Du = A\ (B*u(1:nPt,n)- .5*a^2*dt^2*K*u_bar + f);
107 -             u(1:nPt,n+1) = u(1:nPt,n) + Du;
108 -         end
109 -     else

```


Results obtained-



The simulation has been done for 2 problems for $C=0.5$ and the code seems to be working. To check the stability of the method the following results were obtained by varying the courant number for problem 2 which seems to be consistent with the theory. TG3-2S method is stable for courant number $C^2 \leq 1$, which is consistent with the results. It is unstable for $C > 1$ but stable for values less than equal to 1.



For one dimensional computations both TG3 and TG3-2S are expected to give similar results. However because of the higher dissipative nature of TG3-2S, it is expected to perform better for two-dimensional computations.