# Class Homework 4: 2D Unsteady Convection-Diffusion

By Domingo Eugenio Cattoni Correa:

#### **Two-step fourth-order method:**

Before implementing this method was necessary developed by hand. The Figure 1, shows step by step how the method was obtained by hand.

Two-step Boeth-order netwood
$(t + (g, \nabla)u = S  in this case S = 0$
$\frac{1}{2} u = u + \frac{1}{3} \frac{d}{dt} u_{t} + \frac{1}{12} \frac{dt}{u_{t}t} u_{t} + \frac{u_{t}}{u_{t}} \frac{u_{t}}{u_{t}} u_{t} = \frac{u_{t}}{u_{t}} = \frac{1}{2} \frac{u_{t}}{u_{t}} $
$\frac{2}{2} u^{n+1} = u^n + 4t u^n + 1 4t^2 u^2 = -2 u_{n-1} + 1 = 0 - 0 u^{n+1} + 1 = 0$
12 step compose the second dorivetive of an
$u_t = -(2, 2)u_i  u_{tt} = -(a, 7)^2 u_i$
Now, these two paraveture will be plugged on the method
$\overline{u} = u^{n} + \frac{1}{3} \Delta t (-2 \cdot 2) u^{n} + \frac{1}{12} \Delta t^{2} (-2 \cdot 2)^{n} u^{n}$
Applying Galerich on 1° step
$(w, \mathcal{U}) = (\omega, u^{n}) + \frac{1}{3} \Delta t (w, (-\alpha, \overline{y})u^{n}) + \frac{1}{12} \Delta t^{2} (w, (-\alpha, \overline{y})^{2} u^{n})$
by Parts by parts
(w, (-a, z)w) = -(w, w) + (zw, w)
$ (w_1 (-2.3)^2 u^n) = -a^2 (w, \forall u^n) + (\forall w \forall u^n) $
$(w, \tilde{u}) = (w, w) - \frac{1}{3} 4 \left[ (w, w) - (\nabla w, w) \right] - \frac{1}{12} \Delta t a^2 \left[ a^2 (w, \nabla w) \right] - (\nabla w, \nabla w) \right]$ $= \frac{1}{12} \left[ a^2 (w, \nabla w) \right] - \frac{1}{12} \left[ a^2 (w, \nabla w) \right]$
$MO' = MU'' + \frac{1}{3}\Delta t CU'' - \frac{1}{3}\Delta t M_0 U'' + \frac{1}{12}\Delta t d^2 KU'' - \frac{1}{12}\Delta t d^2 C_0 U''$
Something similar for 2 step. Become of the 1° step
$\alpha^{n+1} = \alpha^n + \Delta t \left(-\alpha \cdot \nabla \omega^n\right) + \frac{1}{2} \Delta t^2 \left(-(\alpha \cdot \nabla)^{2} \widetilde{\omega}\right)$
Applizing Carlierin
$(w, w^{n}) = (w, w^{n}) + \delta t(w, -a, \overline{v}u^{n}) + \frac{1}{2} \delta t^{2} (w, -(2, \overline{v})^{n} t) - \cdot \cdot$
by parts oy Parts
$(w, w^{n+1}) = (w, w^n) - \delta te[(w, w)] - (\nabla w, w)] - \pm \delta t^2 a^2 [(\nabla w, \nabla u) + (w, \nabla u)_{n+1}]$
$MU^{n+1} = MU^{n} + StCU^{n} - \frac{1}{2} \Delta ta^{2} KU^{n} - Sta M_{0}U^{n} - \frac{1}{2} \Delta t^{2}a^{2}C_{0}U^{n}$
Summers: 2 1 Skr > $M \vec{U} = M U^n + 1 \Delta t C U^n - 1 \Delta t M U^n + 1 \Delta t \sigma^2 K U^n - 1 \Delta t \sigma^2 C U^n$
$2^{\circ}$ sho > MU <sup>0H</sup> = MU <sup>0</sup> + $\Delta \pm CU^{\circ}$ - $\Delta \pm c^{2}$ KŨ - $\Delta \pm c$ M <sub>0</sub> U <sup>0</sup> - $1$ $\Delta E c^{2}$ CoU = : C

Figure 1: Method developed by hand.

The next figure shows the line of the code where this method was implemented.



In the code can be observed all the matrices that was developed by hand.

## **Result:**

The problem parameters used in order to show results obtained by using different method were:

## Parameters of the problem:

v(x,y) = (-y,x) (convective velocity)
h = 0.04 (mesh size)
25 bilinear quadrilateral elements (Quad4) per side were used.
S = 0 (source term)

## Initial and boundary condition:

Cosine hill was used as initial condition. Homogeneous Dirichlet condition on the inlet boundary was used.

Eight methods were used in order to computed the solution and see which solution the best is.

#### Galerkin formulation + Lax-Wendroff method with with consistent mass matrix (TG2)



Figure 3: Solution obtained by TG2 a) Isolines, b) 3D plot

Galerkin formulation + Lax-Wendroff method with lumped mass matrix (LW-FD)



Figure 4: Solution obtained by LW-FD a) Isolines, b) 3D plot

## Galerkin formulation + third-order explicit Taylor-Galerkin method (TG3)



Figure 5: Solution obtained by TG3 a) Isolines, b) 3D plot

# Galerkin formulation + Crank-Nicolson with consistent mass matrix (CN)



Figure 6: Solution obtained by CN a) Isolines, b) 3D plot

## Galerkin formulation + Crank-Nicolson with lumped mass matrix (CN-FD)



Figure 7: Solution obtained by CN-FD a) Isolines, b) 3D plot

## Least-squares formulation + Crank-Nicolson method (Carey-Jiang)



Figure 8: Solution obtained by CJ a) Isolines, b) 3D plot

# Galerkin formulation + third order two-step Taylor-Galerkin method (TG3-2S)



Figure 9: Solution obtained by TG3-2S a) Isolines, b) 3D plot

#### Two-step fourth-order (TG4-2S)



Figure 10: Solution obtained by TG4-2S a) Isolines, b) 3D plot

To compare the accuracy of the various methods, explicit and implicit, the maximum and minimum values of the computed solution are provided in Table 1.

Method	Numerical solution		Exact solution		Error %	
	U <sub>max</sub>	U <sub>min</sub>	U <sub>max</sub>	U <sub>min</sub>	$\Delta \text{error\_u}_{\text{max}}$	$\Delta \text{error\_u}_{\min}$
TG2	0.917	-0.0305	1	0	8.3	3.05
LW-FD	0.7340	-0.1947	1	0	27	19.5
TG3	0.9351	-0.0225	1	0	6.5	2.3
CN	0.9855	-0.0541	1	0	1.5	5.4
CN-FD	0.7618	-0.0249	1	0	23.8	2.5
CJ	0.9307	-0.0305	1	0	6.3	3.1
TG3-2S	0.9375	-0.0216	1	0	6.3	2.2
TG4-2S	0.9366	-0.0216	1	0	6.3	2.2

Table 1: Max and min results obtain using different methods.

It can be seen that the worst result obtained was that correspond by using LW-FD.

According to Donae and Huerta's book, it can be split these method into explicit and implicit. The explicit methods used were LW-FD, TG2, TG3, TG3-2S and TG4-2S. The greater accuracy of the finite element schemes employing TG2, TG3, TG3-2S and TG4-2S is clearly apparent, being TG3-2S and TG4+2S the bests methods. On the other hand, the implicit methods used were CN, CJ and CN-FD. CN and CJ gave the best results. It can be observed that the Crank-Nicholson scheme with the Galerking formulation represent a non-dissipative method in pure convection. As a result, significant non-physical oscillation can be seen in Figure 6 and Figure 7. These oscillations were attenuated using a dissipative spatial formulation, such as the least-square FEM of Carey and Jiang. The result can be noted in Figure 8.