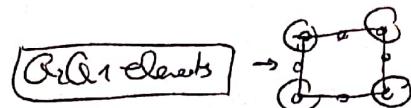


Problem 1

$$\left\{ \begin{array}{l} -V \nabla^2 V + V p = b \text{ in } \Omega \\ \nabla \cdot V = 0 \text{ in } \Omega \\ V = 0 \text{ on } \partial\Omega \end{array} \right. \quad \text{STOKES}$$



a) Q2 Q1 is suitable because it is LBB stable.

For Stokes problem  $(G^{(G^T)})(V_p) = (f_b)$ , we obtain 1<sup>st</sup> equation a substitute in 2<sup>nd</sup>

we obtain:  $V \nabla k^{-1} (f - G^T p) \rightarrow [G K^{-1} G^T] p = G k^{-1} f$

$G K^{-1} G^T$  is SPD if  $\text{ker}(G^T) = \{0\}$ , and this only happens if  $\text{grad } V$  or  $Q$  satisfy LBB condition or inf-sup condition. It is necessary that:

$\dim Q^* \leq \dim V^*$  with  $Q = p \in L_2(\Omega)$   
 $V = V \in H^1(\Omega)$



b) Considering GRS for Stokes & Q2 Q1, we have linear shape pressure basis and quadratic for velocity.

By adding the stabilization:  $\sum_{e \in \partial\Omega} T_e (w, q) \cdot (L(u_p) - F) \Delta e$

we add stabilization parameter to our system:

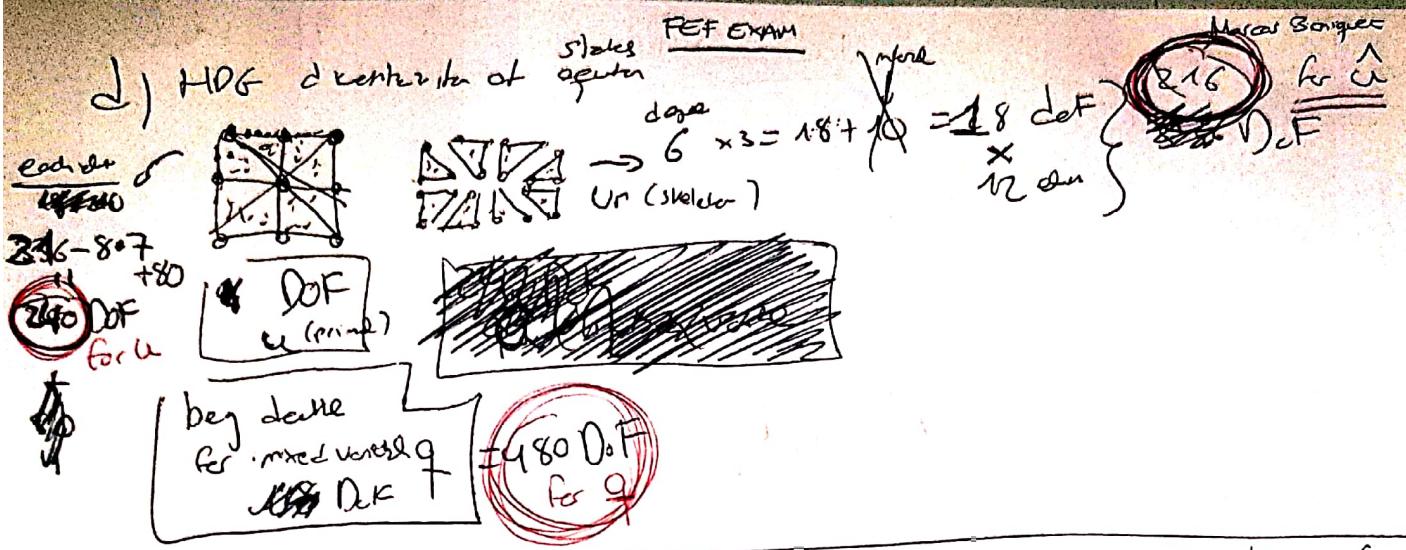
$$\begin{bmatrix} K + K_s & G^T + G \\ G & 0 \end{bmatrix} \begin{bmatrix} V \\ P \end{bmatrix} = \begin{bmatrix} F + S \\ C \end{bmatrix} \quad \text{with } K(G, G)$$

given that velocity is non-linear, all of this additions to stabilize Stokes matters. LBB stable condition was not sufficient.  $T_1, T_2$  stalk be forced (weakly  $\approx 0$ )

c) Possible approach to fulfill LBB condition - degree 6.  $\text{Df.}$

$$\dim Q^* \leq \dim V^* = 6$$

Pressure operator must be of degree  $\leq 6$



e) HDG algorithms: One unknown is a function on  $\Omega$  and the other the trace of this function at the boundaries.



Fig. 2 is dated the global symm ( $\hat{u}$ ) and  $\bar{z}$ , the latter being an auxiliary set of variables.

being an auxiliary set of variables.  
Next thing is to calculate the local problem, calculating primal  
variable  $u$  and mixed variable  $\varphi^0$ .

Optionally: Superconverge may be applied. The converge rate for both  $u$  &  $q$  is  $O(h^{p+1})$ , but a superconvergent algorithm could have up to  $O(h^{p+2})$ . This postprocessed  $\alpha^*$  is calculated solving local problem derived by element.

\* Of course global and local problem must be the first to be established, regarding the mixed variable. Then setting a broken local system of a set of  $\mathcal{L}_2$  subdomains and ~~and~~ setting the primal variable as ~~the~~ be equal to the hybrid variable at the skeleton. For the global problem, transmission conditions on the skeleton and interface conditions are considered.

for the glass f<sub>1</sub>, the  
Nevan Barndy conditions are considered.  
The two sets of conditions are identified within the

Neural Boundary conditions are considered.  
 The numerical fluxes of the HDG model are identified with  
 skeleton one its numerical traces must be stabilized.

Resistive to porosity  $\rightarrow$   $w = \sigma + \frac{\partial u}{\partial t}$  reaction term with  $\sigma > 0$ .

$$\frac{\partial u^m}{\partial t} - V \nabla^2 u^m + (\alpha \cdot \nabla) u^m + \nabla p = f \quad \text{in } \Omega, t > 0$$

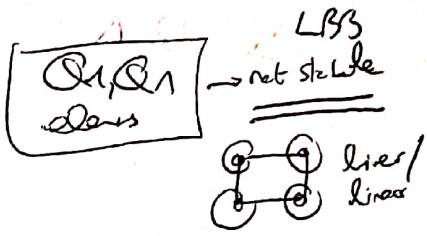
$\nabla \cdot u = 0$

$u = 0$

$u = 0$

new with respect  
No zero states

in  $\Omega, t > 0$   
 $\alpha, \Delta \alpha, \epsilon > 0$   
 $\nabla \cdot \epsilon = 0$



a) Discrete time problem CN.

Finite method:

$$\frac{u^{m+1} - u^m}{\Delta t} - \alpha (\epsilon_{e,m}^{m+1} - \epsilon_e^m) = \epsilon_{e,m}^n$$

with  $\Omega = \mathbb{R}$  for Crack-Necking

$$\frac{z(u^{m+1} - u^m)}{\Delta t} - \epsilon_{e,m}^{m+1} = \epsilon_{e,m}^n$$

$$u_e = V \nabla u - (\alpha \cdot \nabla) u - \sigma u - \nabla p + f$$

Resulting in:  $\frac{2u^{m+1}}{\Delta t} - V \nabla^2 u^{m+1} + (\alpha^{m+1} \cdot V) u^{m+1} + \sigma u^{m+1} + \nabla p^{m+1} =$

$$= \frac{2u^m}{\Delta t} + f^m - V \nabla^2 u^m + (\alpha^m \cdot V) u^m + \sigma u^m + \cancel{\nabla p^m} + \nabla p^n$$

1<sup>st</sup> eqn  $\rightarrow$

$$\frac{z^2}{\Delta t} \epsilon_{e,m}^{m+1} \rightarrow \nabla \cdot u^{m+1} = 0 \quad (\text{incompressibility constraint})$$

b) weak form:

$$\int_{\Omega} w \left[ \left( \frac{z}{\Delta t} + \sigma \right) u^m - V \nabla^2 u^m + (\alpha^m \cdot V) u^m + \nabla p^m \right] \, d\Omega = \int_{\Omega} w \left[ \left( \frac{z}{\Delta t} + \sigma \right) u^m - V \nabla^2 u^m + (\alpha^m \cdot V) u^m + \nabla p^m \right] \, d\Omega$$

$\downarrow$  integrating by parts

$+ \int_{\Gamma_0} w : (\nabla \cdot u^m)$

$\downarrow$  integrating by parts

$\int_{\Gamma_0} w : (\nabla \cdot u^m)$

Considering Dirichlet B.C., with  $u|_{\Gamma_0} = 0$

$$\therefore \int_{\Gamma_0} q^m : (\nabla \cdot u^m) = 0$$

$\leftarrow$  such that  $q^m \in Q^m = L_2(\Omega)$   
 $\leftarrow$  if  $q^m \in Q^m = H^1(\Omega) : u|_{\Gamma_0} = 0$

c)  $\int_{\Omega} \left[ \frac{(\dot{z} + \sigma)}{\Delta t} \left( \underbrace{(\text{mat } N)^T (\text{mat } N)}_{M} \right) \Delta x + V \underbrace{(\text{grad } N)^T (\text{grad } N)}_{K} \Delta x + C(u) \Delta x \right] \Delta x$

$+ \int_{\Omega} \left[ \underbrace{f^T R^T}_{G^T} \right] P = \omega \cdot \left[ \frac{(\dot{z} + \sigma)}{\Delta t} \int_{\Omega} M \cdot u^n + V \int_{\Omega} K \cdot u^n + C(u) \int_{\Omega} u^n \right] + \int_{\Omega} f^T P$

for other form  
 $Q \cdot \int_{\Omega} \underbrace{N^T D}_{P^{nn}} u^n = 0$

$G$

rearranging in a matrix form:

$$\begin{pmatrix} \frac{(\dot{z} + \sigma)}{\Delta t} M + V K + C(u) \\ G^T \end{pmatrix} \begin{pmatrix} u^{nn} \\ P^{nn} \end{pmatrix} = \begin{pmatrix} f^n \\ 0 \end{pmatrix}$$

The difference is the addition of  $\sigma$  term  
with  $M - S$

Berg  $(\text{mat } N) = \begin{pmatrix} n_1 & 0 & \dots & 0 \\ 0 & n_2 & \dots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \dots & n_N \end{pmatrix}$

$(\text{grad } N) = \begin{pmatrix} \frac{\partial n_1}{\partial x_1} & 0 & \dots & 0 \\ 0 & \frac{\partial n_2}{\partial x_1} & \dots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \dots & \frac{\partial n_N}{\partial x_1} \end{pmatrix}$

$D = [1, \dots, 1] (\text{grad } N)$

$M = (\text{mat } N)^T (\text{mat } N) \Delta x$   
 $K = (\text{grad } N)^T (\text{grad } N) \Delta x$   
 $G = \int_{\Omega} \hat{N}^T D \Delta x$

see page 6  
for adding FCS stabilization

d) Given that it is a non-linear system ( $C(u)$ ), it is required a non-linear algorithm, such as Picard or Newton-Raphson. Picard is more robust and does not require the Jacobian, while N-R converges much faster, but requires calculating Jacobian at each iteration, starting with  $u^0 = \underline{u}_0$ . It is more sensitive to initial solution  $u^0$ . (it could diverge fast).

If we choose Picard, then:

$$A(x)x = B(x)$$

$$\text{where } \begin{pmatrix} A(x) \\ B(x) \end{pmatrix} = \begin{pmatrix} \frac{(\dot{z} + \sigma)}{\Delta t} M + V K + C(u) & G^T \\ f^n & 0 \end{pmatrix}$$

Solving linear system at each iteration, starting with  $u^0 = \underline{u}_0$  we obtain  $C(u^0)$  and we can solve linearly to obtain  $u^1$ , etc which is the new  $u^1$ , etc

$$u^0 \rightarrow C(u^0) \rightarrow u^1 \rightarrow C(u^1) \rightarrow u^2 \rightarrow \dots \rightarrow u^k$$

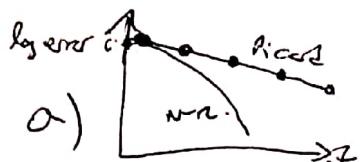
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Problem 2

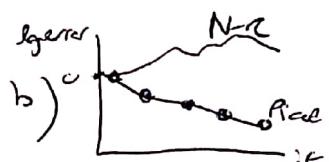
e)  $\delta=0$  (Pure Navier-Stokes problem)

Considering a viscous fluid around object solving [non-dimensional N.S.]

for  $Re=10$  &  $Re=1000$ .



$Re=10$   $\rightarrow$  laminar  
NO turbulence



$Re=1000$   
high turbulence

For an unsteady viscous laminar flow the Ricard regime is unstable, and may appear bifurcations or solution or complex flow patterns. The N.R. case is between this regime and the transition regime which start to have turbulent phenomena.

For Ricard's case ( $Re=10$ ) we have a laminar flow.

- - - - b) ~~so~~ the Jacobian may be affected by the instabilities and or chaotic behaviour on N.R. which may lead to divergence.
- a) Both methods do converge, however N.R. is much faster, as known. Ricard's is an approximation of N.R.

From page 4 FLS [Shear stress] offer unsteady Navier-Stokes with resistance to porosity

 Q1Q1 Not LBB stable. It is necessary to add stabilizer  
 ~~parameters~~ through GLS (Galerkin Least Squares), by imposing that

The weak solution is the solution.

$$\sum_{\text{test func}} \int \nabla \cdot Q(u, q) (L(v, p) - F) dx$$

$$\cancel{\sum_{\text{test func}} \int \left( u \cdot \nabla w + (\omega \cdot \nabla) w \right) + \delta \omega + \nabla q} \left( u \cdot \nabla u + (\omega \cdot \nabla) u + \delta u + \nabla p \right) dx \\ + \sum_{\text{weak eqns}} (\nabla \cdot \omega) (\nabla \cdot u) = 0$$

Important: Q1Q1 is linear for  $p, v$ . Terms with second derivative can be set to 0

$$\begin{aligned} & \cancel{\int \omega \cdot \nabla u} \\ & \cancel{\int \delta u \cdot \nabla u} \\ & \cancel{\int \delta u \cdot \nabla p} \\ & \cancel{\int \delta u \cdot f} \\ & \text{and } \sum_{\text{weak eqns}} (\nabla \cdot \omega) (\nabla \cdot u) = \boxed{K_5} \end{aligned}$$
$$\begin{aligned} & C_1 S(\omega \cdot \nabla) u \\ & C_2 (\omega \cdot \nabla) w (\omega \cdot \nabla) u \\ & C_3 (\omega \cdot \nabla) w \text{ outwards} \\ & C_4 (\omega \cdot \nabla) w \nabla p = \boxed{G_1} \\ & C_5 (\omega \cdot \nabla) w (-f) \\ & C_6 \delta u \nabla p = \boxed{G_2} \\ & C_7 \delta u \nabla f = \boxed{F_w} \\ & C_8 \delta u \nabla \omega = \boxed{G_3} \\ & C_9 \delta u \omega = \boxed{G_4} \\ & C_{10} \delta u = \boxed{G_5} \\ & C_{11} \delta u = \boxed{G_6} \\ & C_{12} \delta u = \boxed{G_7} \\ & C_{13} \delta u = \boxed{G_8} \end{aligned}$$

$$\begin{pmatrix} (\omega, \omega) & (\omega, p) \\ (q, \omega) & (q, p) \end{pmatrix} \begin{pmatrix} u \\ p \end{pmatrix} = \begin{pmatrix} (\omega, f) \\ (q, f) \end{pmatrix}$$

$$\begin{pmatrix} A & G^T \\ G & 0 \end{pmatrix} \begin{pmatrix} u \\ p \end{pmatrix} = \begin{pmatrix} f \\ 0 \end{pmatrix} \quad \text{ori GLS}$$

stabilized:

$$\begin{pmatrix} A + \cancel{K_1} & \cancel{K_2} \\ \cancel{K_3} & G + \cancel{G_1} + \cancel{G_2} \end{pmatrix} \begin{pmatrix} u \\ p \end{pmatrix} = \begin{pmatrix} f + f_w \\ f_w \end{pmatrix}$$