

Ex 1

Inviscid Burgers -  $u_t + uu_x = 0 \quad x \in [0, 1] \quad t > 0$

$$u = u_0(x) \quad x \in [0, 1] \quad t = 0$$

$$u = 1 \quad x = 0 \quad t > 0$$

$$u_0(x) = \begin{cases} 1 & x \in [0, p] \\ 1 - \frac{x-p}{q-p} & x \in [p, q] \\ 0 & x \in (q, 1] \end{cases}$$

$$p = 0.64$$

$$q = 0.84$$

~~Spatial discretization of inviscid burgers of TG-2 approx. is-~~

$$\left( \omega, \frac{\Delta u}{\Delta t} \right) = - \left( \omega, a \cdot \nabla u^n - \frac{\Delta t}{2} (a \cdot \nabla)^2 u^n \right) + \left( \omega, s^n + \right.$$

$$\left. \frac{\Delta t}{2} (s_t^n - a \cdot \nabla s^n) \right)$$

Integrating this by parts becomes,

$$\begin{aligned} \left( \omega, \frac{\Delta u}{\Delta t} \right) &= \left( a \cdot \nabla \omega, u^n + \frac{\Delta t}{2} [s^n - (a \cdot \nabla) u^n] \right) - \\ &\quad \left( (a \cdot n) \omega, u^n + \frac{\Delta t}{2} [s^n - (a \cdot \nabla) u^n] \right)_{\Gamma_{out}} + \\ &\quad (\omega, h^{n+1/2})_{\Gamma_N^{in}} + \left( \omega, s^n + \frac{\Delta t}{2} s_t^n \right) \end{aligned}$$

a)

TG2 approx  $\rightarrow$

$$\frac{u^{n+1} - u^n}{\Delta t} = u_t^n + \frac{\Delta t}{2} u_{tt}^n$$

$$\text{here, } u_t = -f_x$$

$$u_{tt} = -f_{xt} = -f_{tx} = -(f_u u_t)_x = (u^2 u_x)_x.$$

Burgers equ. becomes  $\rightarrow$

$$\frac{u^{n+1} - u^n}{\Delta t} = -f_x^n + \frac{\Delta t}{2} ((u^n)^2 u_x^n)_x$$

②

After integration by parts,

$$\int_0^L w \frac{u^{n+1} - u^n}{\Delta t} dx = \int_0^L w_x f^n dx - \frac{\Delta t}{2} \int_0^L w_x (u^n)^2 u_x^n dx$$

$$\left[ w \left( f^n - \frac{\Delta t}{2} (u^n)^2 u_x^n \right) \right]_{x=0}^{x=L}$$

The boundary term (last term on RHS) can be written as -

$$= - \left[ w \left( f^n - \frac{\Delta t}{2} (u^n)^2 u_x^n \right) \right]_{x=0}^{x=L}$$

$$= - \left[ w \left( f^n + \frac{\Delta t}{2} f_t^n \right) \right]_{x=0}^{x=L}$$

$f$  is element-wise const. representation.

c)

Alternative choices to represent fluxes (non-linear) are -

### ① Classical Representation

$f$  is determined from ' $u$ ' at 2 element gauss points & a two-point Gaussian quadrature is used to evaluate convective term.

### ② Group representation

$f$  is linearly interpolated using its evaluation at the element nodes.

Ex. 2

$$-\nu u_{xx} + \beta u_x = 0 \quad x \in (-1, 1)$$

$$u = 0 \quad x = -1$$

$$u = 1 \quad x = 1$$

$$\nu = 0.03$$

$$\beta = 1.8$$

a)

$$\beta u_x - \nu u_{xx} = 0.$$

The weak form of this equ. is -

$$\int_{-1}^1 (w \cdot \beta u_x + w_x \cdot \nu u_x) dx = \cancel{\int_{-1}^1 w_x (\beta u_x + \nu u_{xx}) dx} = 0$$

$$\text{Shape func. are } N_1(\xi) = \frac{1}{2}(1-\xi), \quad N_2(\xi) = \frac{1}{2}(1+\xi)$$

Discrete equ. at nodes takes the form

$$\int_{-1}^1 \sum \left( \beta N_A \frac{\partial N_B}{\partial x} + \nu \frac{\partial N_A}{\partial x} \frac{\partial N_B}{\partial x} \right) u_B dx = 0$$

The convection & diffusion matrices  $C^e$  &  $K^e$  resp. are -

$$C^e = \beta \int_{-1}^1 \begin{bmatrix} N_A \frac{\partial N_A}{\partial x} & N_A \frac{\partial N_B}{\partial x} \\ N_B \frac{\partial N_A}{\partial x} & N_B \frac{\partial N_B}{\partial x} \end{bmatrix} dx$$

$$K^e = \nu \int_{-1}^1 \begin{bmatrix} \frac{\partial N_A}{\partial x} \frac{\partial N_A}{\partial x} & \frac{\partial N_A}{\partial x} \frac{\partial N_B}{\partial x} \\ \frac{\partial N_B}{\partial x} \frac{\partial N_A}{\partial x} & \frac{\partial N_B}{\partial x} \frac{\partial N_B}{\partial x} \end{bmatrix} dx$$

Discrete form of Galerkin finite element formulation -

$$\beta \left( \frac{u_{j+1} - u_{j-1}}{2h} \right) - \nu \left( \frac{u_{j+1} - 2u_j + u_{j-1}}{h^2} \right) = 0.$$

b)

SGS method equ -

$$\int \omega (\beta \cdot \nabla u) dx + \int \nabla \omega \cdot (\nu \nabla u) dx + \boxed{\sum_e \int P(u) T R(u) dx} = 0$$

$$\hookrightarrow \sum_e \int (\beta \cdot \nabla u + \nabla \cdot (\nu \nabla u)) T (\beta \cdot \nabla u - \nabla \cdot (\nu \nabla u)) dx$$

stabilization term.

Galerkin method is unstable for  $Pe > 1$ . But it gives good results for  $Pe \ll 1$ .

SGS method overcomes the instabilities introduced by GLS method which is supposed to be more stable for  $Pe > 1$ .

$\leftarrow$  Pe in this question is  $> 1$ .  $\therefore$  Galerkin will show oscillations.

$$Pe = \frac{\beta h}{2\nu}$$

SGS method will provide smooth results at  $Pe > 1$ .

c)

Consistent stabilization means ~~more~~ diffusion is added where the mesh is finer. & smooth the sol.

Consistent stabilization gives less numerical diffusion when the numerical sol. appears more close to exact sol.

The stabilization parameter should ~~not~~ vanish as mesh is refined.

Stabilization parameter is influenced by, diffusivity, mesh size, convection vel., reaction cof (if present)