

On "Isoparametric representation"

Assignment 4.1

A 3-node straight bar element is defined by 3 nodes: 1, 2 and 3 with axial coordinates x_1 , x_2 and x_3 respectively as illustrated in figure below. The element has axial rigidity EA , and length $l = x_2 - x_1$. The axial displacement is $u(x)$. The 3 degrees of freedom are the axial node displacement u_1 , u_2 and u_3 . The isoparametric definition of the element is

$$\begin{bmatrix} 1 \\ x \\ u \end{bmatrix} = \begin{bmatrix} 1 & 1 & 1 \\ x_1 & x_2 & x_3 \\ u_1 & u_2 & u_3 \end{bmatrix} \begin{bmatrix} N_1^e \\ N_2^e \\ N_3^e \end{bmatrix} \quad (7.1)$$

in which $N_i^e(\xi)$ are the shape functions of a three bar element. Node 3 lies between 1 and 2 but is not necessarily at the midpoint $x = l/2$. For convenience define,

$$x_1 = 0 \quad x_2 = l \quad x_3 = \left(\frac{l}{2} + \alpha\right)l \quad (7.2)$$

where $-\frac{1}{2} < \alpha < \frac{1}{2}$ characterizes the location of node 3 with respect to the element center. If $\alpha=0$ node 3 is located at the midpoint between 1 and 2.

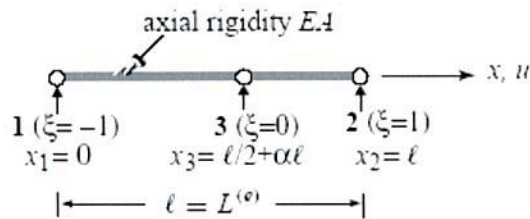


Figure.- The three-node bar element in its local system

1. From (7.2) and the second equation of (7.1) get the Jacobian $J = dx/d\xi$ in terms of l , α and ξ . Show that,

- if $-\frac{1}{4} < \alpha < \frac{1}{4}$ then $J > 0$ over the whole element $-1 \leq \xi \leq 1$.
- if $\alpha = 0$, $J = l/2$ is a constant over the element.

2. Obtain the 1×3 strain displacement matrix \mathbf{B} relating $\mathbf{e} = du/dx = \mathbf{B}\mathbf{u}^e$ where \mathbf{u}^e is the column 3-vector of the node displacement u_1 , u_2 and u_3 . The entries of \mathbf{B} are functions of l , α and ξ .

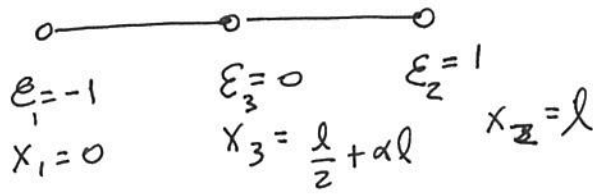
Hint: $\mathbf{B} = d\mathbf{N}/dx = J^{-1}d\mathbf{N}/d\xi$, where $\mathbf{N} = [N_1 \ N_2 \ N_3]$ and J comes from item a).

①

Assignment 4 CSMD

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④.1 Three noded bar



The shape functions for this element can be obtained using.

$$N_j = \prod_{\substack{i=1 \\ i \neq j}}^3 \frac{(\epsilon_i - \epsilon)}{(\epsilon_i - \epsilon_j)} \quad (\text{for 3 nodes})$$

$$N_1 = \frac{(\epsilon_2 - \epsilon)(\epsilon_3 - \epsilon)}{(\epsilon_2 - \epsilon_1)(\epsilon_3 - \epsilon_1)} = \frac{(1 - \epsilon)(0 - \epsilon)}{[1 - (-1)](-(-1))} = \frac{\epsilon}{2}(1 - \epsilon)$$

$\underbrace{\hspace{1.5cm}}_{j=1 \quad i=2} \quad \underbrace{\hspace{1.5cm}}_{j=1 \quad i=3}$

where $\begin{cases} \epsilon_1 = -1 \\ \epsilon_2 = 1 \\ \epsilon_3 = 0 \end{cases}$ similarly for N_2 and N_3

$$N_2 = \frac{(\epsilon_1 - \epsilon)(\epsilon_3 - \epsilon)}{(\epsilon_1 - \epsilon_2)(\epsilon_3 - \epsilon_2)} = \frac{\epsilon}{2}(\epsilon + 1)$$

$$N_3 = \frac{\epsilon_1 - \epsilon}{\epsilon_1 - \epsilon_3} \frac{\epsilon_2 - \epsilon}{\epsilon_2 - \epsilon_3} = 1 - \epsilon^2$$

To represent N_1, N_2 and N_3 in cartesian coordinates we need to find the

Jacobian $\frac{dx}{d\epsilon}$

(2)

since $x = x_1 N_1 + x_2 N_2 + x_3 N_3$

with $x_1 = 0$, $x_2 = l$ & $x_3 = l(\frac{1}{2} + \alpha)$

we get

$$x = 0 + l \frac{\epsilon}{2} (\epsilon + 1) + l(\frac{1}{2} + \alpha)(1 - \epsilon^2)$$

differentiating w.r.t. ϵ

$$\frac{dx}{d\epsilon} = \frac{d}{d\epsilon} \left[\frac{l}{2} (\epsilon^2 + \epsilon) + l(\frac{1}{2} + \alpha) - l(\frac{1}{2} + \alpha)\epsilon^2 \right]$$

$$= \frac{l}{2} (2\epsilon + 1) - 2l(\frac{1}{2} + \alpha)\epsilon$$

$$= \cancel{\epsilon l} + \frac{l}{2} - \cancel{l\epsilon} - 2\alpha\epsilon l$$

$$= \frac{l}{2} - 2\alpha\epsilon l = \boxed{\frac{l}{2} (1 - 4\alpha\epsilon)}$$

the Jacobian $\frac{dx}{d\epsilon} > 0$ if $(1 - 4\alpha\epsilon) > 0$

$$1 > |4\alpha\epsilon|$$

$$\frac{1}{4} > |\alpha||\epsilon|$$

$$-\frac{1}{4} < \alpha < \frac{1}{4}$$

since $|\epsilon| \leq 1$ always

for $\alpha = 0$ the $\frac{dx}{d\epsilon} = \frac{l}{2}$, a constant.

③

Find the strain displacement matrix \bar{B} relating $e = du/dx = \bar{B} \bar{u} e$

$$\left[\bar{B} = \frac{d\bar{N}}{dx} = J^{-1} \frac{d\bar{N}}{dE} \right]$$

Solution

$$B_i = \frac{dN_i}{dx} = \frac{dN_i}{dE} \cdot \frac{dE}{dx} = J^{-1} \frac{dN_i}{dE}$$

using the previous result:

$$J^{-1} = \left[\frac{l}{2} (1 - 4\alpha E) \right]^{-1} = \frac{2}{l(1 - 4\alpha E)}$$

$$\bar{N} = \left[\frac{dN_1}{dE} \quad \frac{dN_2}{dE} \quad \frac{dN_3}{dE} \right]$$

$$= \left[\frac{1}{2} - 2E \quad E + \frac{1}{2} \quad -2E \right]$$

therefore $\bar{B} = J^{-1} \bar{N}$ becomes:

$$\bar{B} = \frac{1}{l(1 - 4\alpha E)} \left[\begin{array}{ccc} 1 - 4E & 2E + 1 & -4E \end{array} \right]$$

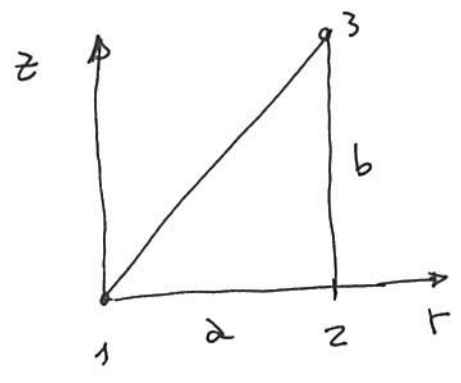
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4.2

compute the entries of the element stiffness matrix K^e for an axisymmetric triangle. Consider $\nu = 0$ (Poisson's ratio) and a stress-strain matrix \bar{D}

$$\bar{D} = E \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1/2 \end{bmatrix}$$

$E = \text{Young's Mod.}$



working with polar coordinates (r, θ, z) symmetry around the z axis

the stiffness matrix is defined as

$$\bar{K} = \int_{\Omega} \bar{B}^T \bar{D} \bar{B} d\Omega$$

to compute the volume of a rotation solid we use

$$K_{ij} = 2\pi \int \bar{B}^T \bar{D} \bar{B} r dr dz$$

where

$$B_i = \begin{bmatrix} \frac{\partial N_i}{\partial r} & 0 \\ 0 & \frac{\partial N_i}{\partial z} \\ \frac{N_i}{r} & 0 \\ \frac{\partial N_i}{\partial z} & \frac{\partial N_i}{\partial r} \end{bmatrix} \quad i = \{1, 2, 3\}$$

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using the definitions of the shape functions

$$W_i = \frac{a_i + b_i r + c_i z}{2A}$$

$$a_i = r_j z_m - r_m z_j$$

$$b_i = z_j - z_m$$

$$c_i = r_m - r_j$$

$$A = \frac{\Delta b}{2}$$

we set
$$N = \frac{1}{\Delta b} \begin{bmatrix} \Delta b - b r & b r - \Delta z & \Delta z \end{bmatrix}$$

$$\frac{\partial N}{\partial r} = \frac{1}{\Delta b} \begin{bmatrix} -b & b & 0 \end{bmatrix}$$

$$\frac{\partial N}{\partial z} = \frac{1}{\Delta b} \begin{bmatrix} 0 & -\Delta & \Delta \end{bmatrix}$$

therefore

$$\bar{\bar{B}} = \frac{1}{\Delta b} \begin{bmatrix} -b & 0 & b & 0 & 0 & 0 \\ 0 & 0 & 0 & -\Delta & 0 & \Delta \\ \frac{\Delta b}{r} - b & 0 & b - \frac{\Delta z}{r} & 0 & \frac{\Delta z}{r} & 0 \\ 0 & -b & -\Delta & b & \Delta & 0 \end{bmatrix}$$

The product $\bar{B}^T \bar{D} \bar{B}$ is dependent on r and z therefore the integration is not trivial.

it can be done element by element

an alternative is to fix r and z at their

centroid values

$$r = \frac{1}{3} \sum_1^3 r_i = \frac{2b}{3}$$

$$z = \frac{1}{3} \sum_1^3 z_i = \frac{b}{3}$$

(b) this way we can evaluate directly

$$K = \int 2\pi \bar{\mathbf{B}}^T \bar{\mathbf{D}} \bar{\mathbf{B}} r dr dz = 2\pi \bar{\mathbf{B}}_c^T \bar{\mathbf{D}}_c \bar{\mathbf{B}}_c r_c A$$

c = centroid.

$$K = \frac{E\pi}{\partial b} \begin{bmatrix} \frac{5b^2}{4} & 0 & -\frac{3b^2}{4} & 0 & \frac{b^2}{4} & 0 \\ \cdot & \frac{b^2}{2} & \frac{\partial b}{2} & -\frac{b^2}{2} & -\frac{\partial b}{2} & 0 \\ \cdot & \cdot & \frac{5b^2 + 2\partial^2}{4} & -\frac{\partial b}{2} & \frac{b^2 - 2\partial^2}{4} & 0 \\ \cdot & \cdot & \cdot & \frac{2\partial^2 + b^2}{2} & \frac{\partial b}{2} & -\partial^2 \\ \text{Symm} & \cdot & \cdot & \cdot & \frac{b^2 + 2\partial^2}{4} & 0 \\ \cdot & \cdot & \cdot & \cdot & \cdot & \partial^2 \end{bmatrix}$$

(2) Rows 2, 4 and 6 add up to zero.

this is the equilibrium in the z direction has to be guaranteed by the stiffness matrix as the problem is not constrained on the other hand, rows 1, 3 and 5 do not sum up to zero. this is because the axial symmetry guarantees equilibrium.

⑦

③ Compute the vector of body forces

$$\bar{f} = -2\pi \int_V N \bar{b} r dr dz$$

again, using the centroid values to avoid integration

$$\bar{f} = -2\pi \bar{N}(r_c, z_c) \begin{bmatrix} b_r \\ b_z \end{bmatrix} r_c A$$

$$\bar{f} = -\frac{2}{9} \pi a^2 b \begin{bmatrix} 1 & 0 \\ 0 & 1 \\ 1 & 0 \\ 0 & 1 \\ 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ g \end{bmatrix} = -\frac{2}{9} \pi a^2 b g \begin{bmatrix} 0 \\ 1 \\ 0 \\ 1 \\ 0 \\ 1 \end{bmatrix}$$