

- Initially square
- sides cannot change their length

1. Deformation mapping  $\varphi$

$$\varphi = \left\{ \begin{array}{l} X_1 + \sin \alpha X_2 \\ \cos \alpha X_2 \end{array} \right\}$$

2.  $\underline{F}$ ,  $\underline{C}$

$$\underline{F} = \frac{d\varphi}{d\underline{x}} = \begin{bmatrix} 1 & \sin \alpha \\ 0 & \cos \alpha \end{bmatrix}$$

$$\underline{C} = \underline{F}^T \underline{F} = \begin{bmatrix} 1 & 0 \\ \sin \alpha & \cos \alpha \end{bmatrix} \begin{bmatrix} 1 & \sin \alpha \\ 0 & \cos \alpha \end{bmatrix} = \begin{bmatrix} 1 & \sin \alpha \\ \sin \alpha & 1 \end{bmatrix}$$

3. Variation in volume

$$dv = J dV \quad J = \det F$$

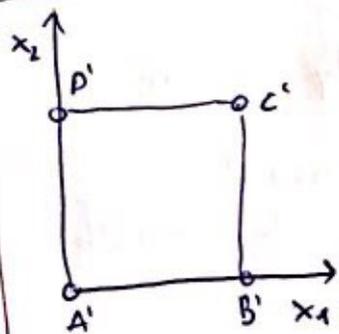
$$\det \underline{F} = \cos \alpha \quad \rightarrow \quad \boxed{dv = \cos \alpha dV}$$

$$\alpha = 0$$

$$\alpha = \pi/4$$

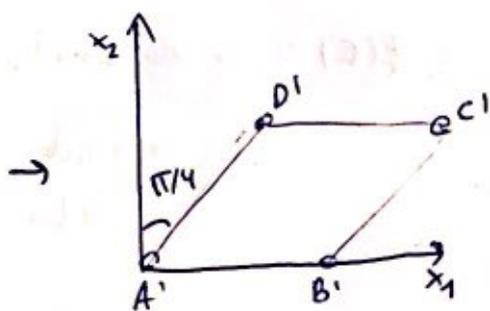
$$\alpha = \pi/2$$

\*

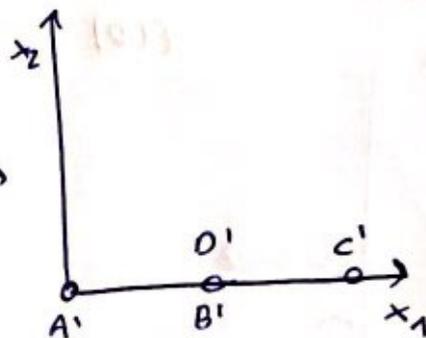


$$\frac{dv}{dV} = 1$$

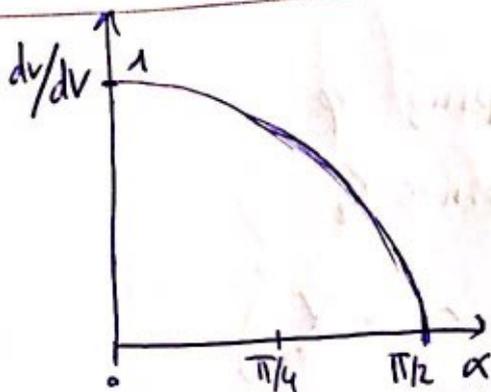
(No change)



$$\frac{dv}{dV} = \frac{\sqrt{2}}{2} \approx 0.7$$



$$\frac{dv}{dV} = 0$$



The volume continuously decreases as  $\alpha$  increases.

4. When deformations are not admissible

If  $J \leq 0$  the deformation is not admissible

In this case, as  $J = \cos \alpha$ , for

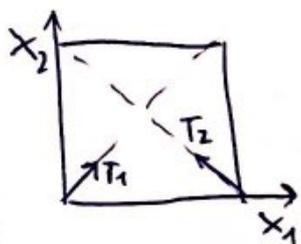
$\alpha \geq \pi/2$  the deformation is not admissible.

\* See the geometric interpretation above

5. de og diagonals AC, BD

dp being  $\beta$   $\widehat{AC BD}$

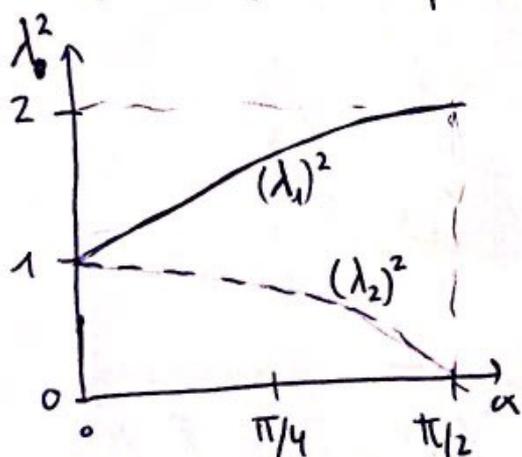
stretching:  $\lambda^2 = T \cdot C \cdot T$  being  $T$   $\begin{cases} T_1 = \frac{1}{\sqrt{2}}(1, 1) \\ T_2 = \frac{1}{\sqrt{2}}(-1, 1) \end{cases}$



$$\lambda_1^2 = \frac{1}{2}(1, 1) \begin{pmatrix} 1 & \sin \alpha \\ \sin \alpha & 1 \end{pmatrix} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = 1 + \sin \alpha$$

$$\lambda_2^2 = \frac{1}{2}(-1, 1) \begin{pmatrix} 1 & \sin \alpha \\ \sin \alpha & 1 \end{pmatrix} \begin{pmatrix} -1 \\ 1 \end{pmatrix} = 1 - \sin \alpha$$

therefore, if we plot the expressions  $\lambda_1(\alpha)$ ,  $\lambda_2(\alpha)$



Here, we can see that the diagonal AC will reach a maximum when  $\alpha = \pi/2$  (4 points on the horizontal).

Meanwhile, diagonal BD will ~~be~~ decrease

until reaches its minimum at  $\alpha = \pi/2$

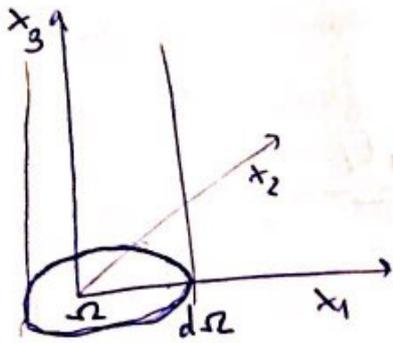
Angle

$$\cos \beta = \frac{T_1 \cdot C \cdot T_2}{\underbrace{\sqrt{T_1 \cdot C \cdot T_1}}_{\lambda_1} \underbrace{\sqrt{T_2 \cdot C \cdot T_2}}_{\lambda_2}}$$

~~Be~~  $\cos \beta = 0$  ( $90^\circ$ )  
constant

$$T_1 \cdot C \cdot T_2 = \frac{1}{2}(1, 1) \begin{pmatrix} 1 & \sin \alpha \\ \sin \alpha & 1 \end{pmatrix} \begin{pmatrix} -1 \\ 1 \end{pmatrix} = -1 - \sin \alpha + 1 + \sin \alpha = 0$$

## HOMEWORK B C



Anti-plane shear deformation:

$$\begin{cases} \varphi_1 = x_1 \\ \varphi_2 = x_2 \\ \varphi_3 = x_3 + w(x_1, x_2) \end{cases}$$

① sketch deformation of  $\Omega$ 

$$a. \quad \underline{F} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ \frac{dw}{dx_1} & \frac{dw}{dx_2} & 1 \end{pmatrix} \quad \underline{C} = \underline{F}^T \underline{F}$$

$$\underline{C} = \begin{pmatrix} 1 + \left(\frac{dw}{dx_1}\right)^2 & \frac{dw}{dx_1} \frac{dw}{dx_2} & \frac{dw}{dx_1} \\ \frac{dw}{dx_1} \frac{dw}{dx_2} & 1 + \left(\frac{dw}{dx_2}\right)^2 & \frac{dw}{dx_2} \\ \frac{dw}{dx_1} & \frac{dw}{dx_2} & 1 \end{pmatrix}$$

Symm.

Jacobian:  $J = \det F$ 

$$J = \begin{vmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ \frac{dw}{dx_1} & \frac{dw}{dx_2} & 1 \end{vmatrix} = 1$$

② b. The volume does not change because

$$dv = J dV \quad \text{and} \quad J = 1$$

c. The local impenetrability condition is satisfied because  $\det F > 0$ .

## 2. Unit vectors

$$A = \frac{w_1 E_1 + w_2 E_2}{\sqrt{w_1^2 + w_2^2}}, \quad B = \frac{-w_2 E_1 + w_1 E_2}{\sqrt{w_1^2 + w_2^2}}$$

a.

Assuming that  $w(x_1, x_2)$  is a contour level, its gradient is:  $\nabla w = w_1 E_1 + w_2 E_2$ .

So,  $A$  is a normal vector over the contour level

$$A = \frac{\nabla w}{\|\nabla w\|}$$

Given that  $A \cdot B = 0 \rightarrow B$  is perpendicular and, therefore, is a vector contained on  $w(x_1, x_2)$ .

b. change in length and angle

$$\begin{aligned} \frac{da}{dA} &= \sqrt{T_A \cdot \underline{\underline{c}} \cdot T_A} \rightarrow \text{unit vector} \\ &= \sqrt{\left(\frac{dw}{dx_1}\right)^2 + \left(\frac{dw}{dx_2}\right)^2 + 1} \quad \circ \end{aligned}$$

\* calculations done with the MATLAB symbolic tool

$$\frac{db}{dB} = \sqrt{T_B \cdot \underline{\underline{c}} \cdot T_B} = 1$$

For the angle:

$$\cos(\underline{A}, \underline{B}) = \frac{\sqrt{T_A \cdot \underline{\underline{c}} \cdot T_B}}{\frac{da}{dA} \frac{db}{dB}}$$

this term is equal to zero.  $\theta = 90^\circ$

conclusions:

A and B remains orthogonal in the deformed configuration.

Regarding the length, B remains constant and A change as a function of  $w(x_1, x_2)$ .

The results coincides with the definition of shear antiplane shear, where the body displacements are 0 (B) but are not null in the perpendicular direction (A).

### 3. Piola transformation

$$d\underline{s} = \underline{J} \underline{F}^{-T} d\underline{S}$$

↳ inverse of transpose of  $F$

~~the reference unit vector of  $d\underline{S}$ :~~ the reference unit vector of  $d\underline{S}$  =

$$T_{dS} = [0, 0, 1]$$

$$d\underline{s} = 1 \cdot \begin{bmatrix} 1 & 0 & -dw/dx_1 \\ 0 & 1 & -dw/dx_2 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} dS = \begin{bmatrix} -\frac{dw}{dx_1} \\ -\frac{dw}{dx_2} \\ 1 \end{bmatrix} dS$$

5. Deformed area of  $\Omega$

From the previous expression:

$$d\underline{s} = \underline{J} \underline{F}^{-T} d\underline{s} \rightarrow \int_{\Omega^0} d\underline{s} = \int_{\Omega} \|\underline{J} \underline{F}^{-T} d\underline{s}\| d\Omega$$

$$\int_{\Omega^0} d\underline{s} = \int_{\Omega^0} \sqrt{\left(\frac{dw}{dx_1}\right)^2 + \left(\frac{dw}{dx_2}\right)^2} + 1 d\underline{s}$$

6. boundary  $\partial\Omega$ :

$$x_1^d = x_1(s) \quad x_2 = x_2(s)$$

Being the boundary  $\partial\Omega$

$$\begin{cases} x_1 = x_1(s) \\ x_2 = x_2(s) \end{cases} \quad 0 \leq s \leq L \quad \text{arc-length}$$

Perimeter:  $\varphi$

tangent vector.

$$\varphi = \int_0^L \sqrt{\left(\frac{dx_1(s)}{ds}\right)^2 + \left(\frac{dx_2(s)}{ds}\right)^2} ds \quad \underline{t} = \frac{E_1 x_1'(s) + E_2 x_2'(s)}{\|E_1 x_1'(s) + E_2 x_2'(s)\|}$$

$\lambda_T \rightarrow$  stretching in the direction of  $\underline{t}$

$$\varphi = \int_{\partial\Omega} \lambda_T ds = \int_0^L \|\underline{t} \cdot \underline{\underline{c}} \cdot \underline{t}\| ds$$