# FINITE ELEMENTS 

## Master of Science in Computational Mechanics/ Numerical Methods Fall Semester 2015

Homework 2: Plane Elasticity
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Exercise 1- Describe the strong form of the problem in the reduced domain (left half). Indicate accurately the Boundary Conditions in every edge.

The governing/field equations for linear elasticity are:

- Balance equation : $\nabla \cdot \boldsymbol{\sigma}+\boldsymbol{b}=\rho \ddot{\boldsymbol{u}}$.
- Constitutive law : $\boldsymbol{\sigma}=\boldsymbol{C}: \boldsymbol{\varepsilon} \rightarrow \boldsymbol{\sigma}=\boldsymbol{D} \boldsymbol{\varepsilon}$.
- Kinematic equation: $\boldsymbol{\varepsilon}=\frac{1}{2}\left(\nabla \boldsymbol{u}+(\nabla \boldsymbol{u})^{T}\right)$.

Where:

- $\boldsymbol{\sigma}$ is the Cauchy stress tensor.
- $\boldsymbol{b}$ is the body foces per unit of volume.
- $\quad \rho$ is the density.
- $\boldsymbol{u}$ is the displacement field.
- $\boldsymbol{C}$ is the stiffness tensor of fourth order, which is simplied to a second order $\boldsymbol{D}$.
- $\boldsymbol{\varepsilon}$ is the strain tensor.

For the boundary conditions we have to take into account that the symmetry plane prevents any displacement on the x direction among it, as well as the fact there is no resultant force perpendicular to the plane.

Finally the strong form is described by:

$$
\left\{\begin{array}{l}
\text { Balance equation }: \nabla \cdot \boldsymbol{\sigma}+\boldsymbol{b}=\rho \ddot{\boldsymbol{u}} \\
\text { Dirichlet } B C:\left\{\begin{array}{c}
u=0, v=0 \text { on } x \in[-3,0] \text { and } y=0 \\
u=0 \text { on } x=0 y \in[0,3) \\
u=0, v=-\delta \text { on } x=0 \text { and } y=3
\end{array}\right. \\
\text { Neuman } B C:\left\{\begin{array}{c}
\boldsymbol{\sigma} \cdot \boldsymbol{n}=0 \text { on } x \in[-3,0] \text { and } y=3+x \\
\boldsymbol{\sigma} \cdot \boldsymbol{e}_{x}=0 \text { on } x=0 \text { and } y \in[0,3)
\end{array}\right.
\end{array}\right.
$$

Exercise 2- Describe the mesh shown in figure 2 by giving the arrays of nodal coordinates $\mathbf{X}$ and the connectivity matrix $\mathbf{T}$. In order to simplify the computations select the local numbering of the nodes such that, in every element, the node in the right angle vertex has loca numer equal to 1 .

The vector of nodal coordinates describes the location of the nodes in global numbering. For the case of study:

$$
\boldsymbol{X}=\left[\begin{array}{cc}
x & y \\
-3 & 0 \\
-1.5 & 0 \\
0 & 0 \\
-1.5 & 1.5 \\
0 & 1.5 \\
0 & 3
\end{array}\right] \begin{gathered}
1 \\
2 \\
3 \\
4 \\
5 \\
6
\end{gathered}
$$

where the rows relate with the global numbering of the nodes, and the columns with the $X$ and Y coordinates.

The connectivity matrix is the one which describes the nodes composing an element and also relates the global and local numberings.

Setting the local numbering in counterclockwise direction starting from the right angle, the connectivity matrix is:

$$
\mathrm{T}=\left[\begin{array}{llll}
(1) & (2) & 3 & 4 \\
2 & 4 & 3 & 5 \\
4 & 2 & 5 & 6 \\
1 & 5 & 2 & 4
\end{array}\right] \quad \begin{gathered}
\mathrm{i} \\
\mathrm{j} \\
\mathrm{k}
\end{gathered}
$$

where the columns relate with the elements and the rows with the local numbering of those elements.

Exercise 3. Set up the linear system of equations corresponding to the discretization in figure 2. How many degrees of freedom has the system to be solved?

Parting from (1.a) neglecting all the dependencies on $z$ and time, it results the system:

$$
\left\{\begin{array}{l}
\frac{\partial \sigma_{x}}{\partial x}+\frac{\partial \sigma_{x y}}{\partial y}+b_{x}=0  \tag{2}\\
\frac{\partial \sigma_{x y}}{\partial y}+\frac{\partial \sigma_{y}}{\partial x}+b_{y}=0
\end{array}\right.
$$

Multiplying by test functions $w_{1}$ and $w_{2}$ and integrating over all the domain:

$$
\left\{\begin{array}{l}
\int_{\Omega} w_{1} \frac{\partial \sigma_{x}}{\partial x}+w_{1} \frac{\partial \sigma_{x y}}{\partial y}+w_{1} b_{x} d \Omega  \tag{3}\\
\int_{\Omega} w_{2} \frac{\partial \sigma_{x y}}{\partial y}+w_{2} \frac{\partial \sigma_{y}}{\partial x}+w_{2} b_{y} d \Omega
\end{array}\right.
$$

Splitting the integrals and integrating by parts the two first terms of both equations:

$$
\left\{\begin{array}{l}
-\int_{\Omega} \frac{\partial w_{1}}{\partial x} \sigma_{x} d \Omega+\oint_{\Gamma_{q}} w_{1} \sigma_{x} d \Gamma-\int_{\Omega} \frac{\partial w_{1}}{\partial y} \sigma_{x y} d \Omega+\oint_{\Gamma_{q}} w_{1} \sigma_{x y} d \Gamma+\int_{\Omega} w_{1} b_{x} d \Omega  \tag{4}\\
-\int_{\Omega} \frac{\partial w_{2}}{\partial x} \sigma_{x y} d \Omega+\oint_{\Gamma_{q}} w_{2} \sigma_{x y} d \Gamma-\int_{\Omega} \frac{\partial w_{2}}{\partial x} \sigma_{y} d \Omega+\oint_{\Gamma_{q}} w_{2} \sigma_{y} d \Gamma+\int_{\Omega} w_{2} b_{y} d \Omega
\end{array}\right.
$$

Rearranging, putting it in matrix form and considering the traction vector as $\boldsymbol{t}=\boldsymbol{\sigma} \cdot \boldsymbol{n}$ :

$$
\int_{\Omega}\left[\begin{array}{ccc}
\frac{\partial w_{1}}{\partial x} & 0 & \frac{\partial w_{1}}{\partial y}  \tag{5}\\
0 & \frac{\partial w_{2}}{\partial y} & \frac{\partial w_{2}}{\partial x}
\end{array}\right]\left[\begin{array}{c}
\sigma_{x} \\
\sigma_{y} \\
\sigma_{x y}
\end{array}\right] d \Omega=\int_{\Omega}\left[\begin{array}{c}
w_{1} b_{x} \\
w_{2} b_{y}
\end{array}\right] d \Omega+\oint_{\Gamma_{q}}\left[\begin{array}{l}
w_{1} t_{x} \\
w_{2} t_{y}
\end{array}\right] d \Gamma
$$

Substituting the constitutive law (1.b) in Voigt notation and considering plane stress:

$$
\begin{gather*}
\int_{\Omega}\left[\begin{array}{ccc}
\frac{\partial w_{1}}{\partial x} & 0 & \frac{\partial w_{1}}{\partial y} \\
0 & \frac{\partial w_{2}}{\partial y} & \frac{\partial w_{2}}{\partial x}
\end{array}\right]\left[\begin{array}{ccc}
\frac{E}{1-v^{2}} & \frac{v E}{1-v^{2}} & 0 \\
\frac{v E}{1-v^{2}} & \frac{E}{1-v^{2}} & 0 \\
0 & 0 & \frac{E}{2(1+v)}
\end{array}\right]\left[\begin{array}{c}
\varepsilon_{x} \\
\varepsilon_{y} \\
\gamma_{x y}
\end{array}\right] d \Omega  \tag{6}\\
\\
=\int_{\Omega}\left[\begin{array}{c}
w_{1} b_{x} \\
w_{2} b_{y}
\end{array}\right] d \Omega+\oint_{\Gamma_{q}}\left[\begin{array}{l}
w_{1} t_{x} \\
w_{2} t_{y}
\end{array}\right] d \Gamma
\end{gather*}
$$

Where $E$ is the young modulus and $v$ the Poisson ratio.

For a 3-noded triangle the interpolation for $u^{h}$ and $v^{h}$ is:

$$
\begin{gather*}
u^{h}=N_{i} u_{i}=N_{1} u_{1}+N_{2} u_{2}+N_{3} u_{3}  \tag{7.a}\\
v^{h}=N_{i} v_{i}=N_{1} v_{1}+N_{2} v_{2}+N_{3} v_{3} \tag{7.b}
\end{gather*}
$$

Where:

$$
\begin{array}{ll}
- & N_{i}=\frac{1}{2 A}\left(a_{i}+b_{i} x+c_{i} y\right) \\
- & a_{i}=x_{j} y_{k}-x_{k} y_{j} ; \quad b_{i}=y_{j}-y_{k} ; c_{i}=x_{k}-x_{j}
\end{array}
$$

with A being the area of the element.
From the kinematic equation (1.c):

$$
\begin{align*}
&- \varepsilon_{x}=\frac{\partial u}{\partial x}=\xrightarrow{u^{h} \approx u}=\frac{\partial N_{i} u_{i}}{\partial x}  \tag{8.a}\\
&-\quad \varepsilon_{y}=\frac{\partial v}{\partial y}=\xrightarrow{v^{h} \approx v}=\frac{\partial N_{i} v_{i}}{\partial y}  \tag{8.b}\\
&-\quad \gamma_{x y}=\frac{\partial u}{\partial y}+\frac{\partial v}{\partial x}=\xrightarrow{u^{h} \approx u ; v^{h} \approx v}=\frac{\partial N_{i} u_{i}}{\partial y}+\frac{\partial N_{i} v_{i}}{\partial x} \tag{8.c}
\end{align*}
$$

Substituting (8.a-8.c) into (6) and taking the test functions $w_{i}$ the same as the shape functions $N_{i}$ (Galerkin method):

$$
\left.\begin{array}{l}
\int_{\Omega}\left[\begin{array}{ccc}
\frac{\partial N_{1}}{\partial x} & 0 & \frac{\partial N_{1}}{\partial y} \\
0 & \frac{\partial N_{1}}{\partial y} & \frac{\partial N_{1}}{\partial x} \\
\frac{\partial N_{2}}{\partial x} & 0 & \frac{\partial N_{2}}{\partial y} \\
0 & \frac{\partial N_{2}}{\partial y} & \frac{\partial N_{2}}{\partial x} \\
\frac{\partial N_{3}}{\partial x} & 0 & \frac{\partial N_{3}}{\partial y} \\
\frac{E}{1-v^{2}} & \frac{v E}{1-v^{2}} & 0 \\
\frac{v E}{1-v^{2}} & \frac{E}{1-v^{2}} & 0 \\
0 & 0 & \frac{E}{2(1+v)}
\end{array}\right]\left[\begin{array}{ccccc}
\frac{\partial N_{1}}{\partial x} & 0 & \frac{\partial N_{2}}{\partial x} & 0 & \frac{\partial N_{3}}{\partial x} \\
0 & 0 \\
0 & \frac{\partial N_{1}}{\partial y} & 0 & \frac{\partial N_{2}}{\partial y} & 0 \\
\frac{\partial N_{3}}{\partial x}
\end{array}\right] \\
\frac{\partial N_{3}}{\partial y} \\
\frac{\partial N_{1}}{\partial x} \\
\frac{\partial N_{2}}{\partial y} \\
\frac{\partial N_{2}}{\partial x} \\
\frac{\partial N_{3}}{\partial y} \\
\frac{\partial N_{3}}{\partial x}
\end{array}\right]\left[\begin{array}{l}
u_{1} \\
v_{1} \\
u_{2} \\
v_{2} \\
u_{3} \\
v_{3}
\end{array}\right] d \Omega
$$

Using index notation:

$$
\begin{align*}
& \int_{\Omega} \frac{1}{2 A}\left[\begin{array}{ccc}
b_{i} & 0 & c_{i} \\
0 & c_{i} & b_{i}
\end{array}\right]\left[\begin{array}{ccc}
\frac{E}{1-v^{2}} & \frac{v E}{1-v^{2}} & 0 \\
\frac{v E}{1-v^{2}} & \frac{E}{1-v^{2}} & 0 \\
0 & 0 & \frac{E}{2(1+v)}
\end{array}\right] \frac{1}{2 A}\left[\begin{array}{cc}
b_{j} & 0 \\
0 & c_{j} \\
c_{j} & b_{j}
\end{array}\right] d \Omega a_{i}  \tag{9}\\
&=\int_{\Omega}\left[\begin{array}{c}
N_{i} b_{x} \\
N_{i} b_{y}
\end{array}\right] d \Omega+\oint_{\Gamma_{q}}\left[\begin{array}{l}
N_{i} t_{x} \\
N_{i} t_{y}
\end{array}\right] d \Gamma
\end{align*}
$$

This last expression can be seen as a linear system of equations of the form $\mathbf{K}^{\ominus} \mathbf{a}^{\ominus}=\mathbf{f}{ }^{\ominus}$, where

$$
k_{i j}^{\odot}=\int_{\Omega} \frac{1}{2 A^{\odot}}\left[\begin{array}{ccc}
b_{i} & 0 & c_{i}  \tag{10}\\
0 & c_{i} & b_{i}
\end{array}\right]^{\odot}\left[\begin{array}{ccc}
\frac{E}{1-v^{2}} & \frac{v E}{1-v^{2}} & 0 \\
\frac{v E}{1-v^{2}} & \frac{E}{1-v^{2}} & 0 \\
0 & 0 & \frac{E}{2(1+v)}
\end{array}\right]^{\odot} \frac{1}{2 A^{\odot}}\left[\begin{array}{cc}
b_{j} & 0 \\
0 & c_{j} \\
c_{j} & b_{j}
\end{array}\right]^{\odot} d \Omega
$$

and

$$
f_{i}^{\odot}=\int_{\Omega}\left[\begin{array}{l}
N_{i} b_{x}  \tag{11}\\
N_{i} b_{y}
\end{array}\right]^{\odot} d \Omega+\oint_{\Gamma_{q}}\left[\begin{array}{l}
N_{i} t_{x} \\
N_{i} t_{y}
\end{array}\right]^{\odot} d \Gamma
$$

For the specific case of study we can integrate taking into consideration that all terms inside $k_{i j}^{\ominus}$ are constant, and that the body force $(\boldsymbol{b}=\rho \boldsymbol{g})$ is uniform among the elements:

$$
\begin{align*}
& k_{i j}^{\odot}=\frac{1}{4 A^{\odot}}\left[\begin{array}{ccc}
b_{i} & 0 & c_{i} \\
0 & c_{i} & b_{i}
\end{array}\right]^{\odot}\left[\begin{array}{ccc}
\frac{E}{1-v^{2}} & \frac{v E}{1-v^{2}} & 0 \\
\frac{v E}{1-v^{2}} & \frac{E}{1-v^{2}} & 0 \\
0 & 0 & \frac{E}{2(1+v)}
\end{array}\right]^{\odot}\left[\begin{array}{cc}
b_{j} & 0 \\
0 & c_{j} \\
c_{j} & b_{j}
\end{array}\right]^{\odot}  \tag{12}\\
& f_{i}^{\ominus}=\left[\begin{array}{c}
0 \\
-\rho g
\end{array}\right]+\oint_{\Gamma_{q}}\left[\begin{array}{l}
N_{i} t_{x} \\
N_{i} t_{y}
\end{array}\right]^{\ominus} d \Gamma \tag{13}
\end{align*}
$$

The expressions above refer to the elements and therefore they need to be assembled using the connectivity matrix which relates global and nodal numbering.
For element 1 :

For element 2:

$$
k^{(2)}=\left[\begin{array}{ccc}
4 & 2 & 5 \\
k_{11}^{(2)} & k_{12}^{(2)} & k_{13}^{(2)} \\
k_{21}^{(2)} & k_{22}^{(2)} & k_{23}^{(2)} \\
k_{31}^{(2)} & k_{32}^{(2)} & k_{33}^{(2)}
\end{array}\right] \begin{aligned}
& 4 \\
& 2 \\
& 5
\end{aligned} ; f^{(2)}=\left[\begin{array}{c}
f_{1}^{(2)} \\
f_{2}^{(2)} \\
f_{3}^{(2)}
\end{array}\right] \begin{aligned}
& 4 \\
& 2 \\
& 5
\end{aligned}
$$

For element 3:

$$
\left.k^{(3}=\left[\begin{array}{ccc}
3 & 5 & 2 \\
k_{11}^{(3} & k_{12}^{(3} & k_{13}^{(3} \\
k_{21}^{\Theta} & k_{22}^{\Theta} & k_{23}^{\Theta} \\
k_{31}^{\Theta} & k_{32}^{\Theta} & k_{33}^{\Theta}
\end{array}\right] \quad \begin{array}{c}
3 \\
5
\end{array} ; f^{(3}=\left[\begin{array}{c}
f_{1}^{(3} \\
2
\end{array}\right] \begin{array}{l}
3 \\
f_{2}^{(3)} \\
f_{3}^{(3}
\end{array}\right] \begin{aligned}
& \\
& 2
\end{aligned}
$$

For element 4:

Where the $k_{i j}^{\odot}$ and $f_{i}^{\odot}$ components are to be computed using (10) and (11).
Finally assembling:

$$
\boldsymbol{f}=\left[\begin{array}{c}
f_{3}{ }^{(1)} \\
f_{1}{ }^{(1)}+f_{2}{ }^{(2)}+f_{3}^{(3)} \\
f_{1}^{(3)} \\
f_{2}{ }^{(1)}+f_{1}{ }^{(2)}+f_{3}^{(4)} \\
f_{3}^{(2)}+f_{2}{ }^{(3)}+f_{1}^{(4)} \\
f_{2}{ }^{(4)}
\end{array}\right]
$$

Being the vector of unknowns $\boldsymbol{a}$ :

$$
\boldsymbol{a}=\left[\begin{array}{llllll}
a_{1} & a_{2} & a_{3} & a_{4} & a_{5} & a_{6}
\end{array}\right]^{T} \text {; with } a_{i}=\left[\begin{array}{l}
u_{i} \\
v_{i}
\end{array}\right]
$$

In the case of study there are 6 nodes with 2 degrees of freedom each, totaling 12 degrees of freedom. However after applying the Dirichlet Boundary conditions there remain 3 degrees of freedom $\left(u_{4}, v_{4}\right.$ and $\left.v_{5}\right)$.

Exercise 4. Compute the FE approximation $u^{h}$. Use $E=10 G P a, v=0.2, \delta=10^{-2}$ and $\rho g=10^{3} \mathrm{~N} / \mathrm{m}^{2}$
The assembled stiffness matrix $\boldsymbol{K}$ and force vector $\boldsymbol{f}$ take the form:
$\boldsymbol{K}=10^{10}\left(\begin{array}{ccccccccccccc}0.5208 & 0 & 0.5208 & 0.1042 & 0 & 0 & 0 & 0.1042 & 0 & 0 & 0 & 0 \\ 0 & 0.2083 & 0.2083 & -0.2083 & 0 & 0 & -0.2083 & 0 & 0 & 0 & 0 & 0 \\ -0.5208 & 0.2083 & 14.583 & -0.3125 & -0.5208 & 0.1042 & -0.4167 & 0.3125 & 0 & -0.3125 & 0 & 0 \\ 0.1042 & -0.2083 & -0.3125 & 14.583 & 0.2083 & -0.2083 & 0.3125 & -10.417 & -0.3125 & 0 & 0 & 0 \\ 0 & 0 & -0.5208 & 0.2083 & 0.7292 & -0.3125 & 0 & 0 & -0.2083 & 0.1042 & 0 & 0 \\ 0 & 0 & 0.1042 & -0.2083 & -0.3125 & 0.7292 & 0 & 0 & 0.2083 & -0.5208 & 0 & 0 \\ 0 & -0.2083 & -0.4167 & 0.3125 & 0 & 0 & 14.583 & -0.3125 & -10.417 & 0.3125 & 0 & -0.1042 \\ -0.1042 & 0 & 0.3125 & -10.417 & 0 & 0 & -0.3125 & 14.583 & 0.3125 & -0.4167 & -0.2083 & 0 \\ 0 & 0 & 0 & -0.3125 & -0.2083 & 0.2083 & -10.417 & 0.3125 & 14.583 & -0.3125 & -0.2083 & 0.1042 \\ 0 & 0 & -0.3125 & 0 & 0.1042 & -0.5208 & 0.3125 & -0.4167 & -0.3125 & 14.583 & 0.2083 & -0.5208 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & -0.2083 & -0.2083 & 0.2083 & 0.2083 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -0.1042 & 0 & 0.1042 & -0.5208 & 0 & 0.5208 \\ 0 & & & & & & 0 & 0 & 0 & 0 & 0 & 0 & 0\end{array}\right)$

$$
\boldsymbol{f}=\left(\begin{array}{c}
R_{1 x} \\
R_{1 y}-375 \\
R_{2 x} \\
R_{2 y}-1125 \\
R_{3 x} \\
R_{3 y}-375 \\
0 \\
-1125 \\
R_{5 x} \\
-1125 \\
R_{6 x} \\
R_{6 y}-375 \\
\end{array}\right)
$$

And the displacement vector $\boldsymbol{a}$ with the Dirichlet BC remains:

$$
\boldsymbol{a}=\left[\begin{array}{llllllllllll}
0 & 0 & 0 & 0 & 0 & 0 & u_{4} & v_{4} & 0 & v_{5} & 0 & -\delta
\end{array}\right]^{T}
$$

The only rows with unknowns are 7, 8, and 10; therefore reducing the system:

$$
10^{10}\left[\begin{array}{ccc}
1.4583 & -0.3125 & 0.3125 \\
-0.3125 & 1.4583 & -0.4167 \\
0.3125 & -0.4167 & 1.4583
\end{array}\right]\left[\begin{array}{l}
u_{4} \\
v_{4} \\
v_{5}
\end{array}\right]=\left[\begin{array}{c}
0 \\
-1125 \\
-1125
\end{array}\right]-\left(-10^{-2}\right)\left[\begin{array}{c}
-0.1042 \\
0 \\
-0.5208
\end{array}\right]
$$

Solving:

$$
\left[\begin{array}{l}
u_{4} \\
v_{4} \\
v_{5}
\end{array}\right]=\left[\begin{array}{l}
-1.28 \cdot 10^{-4} \\
-1.13 \cdot 10^{-3} \\
-3.86 \cdot 10^{-3}
\end{array}\right] m
$$

The reactions can be computed 'a posteriori':

$$
\boldsymbol{R}=10^{3}\left[\begin{array}{c}
1179 \\
267 \\
9081 \\
11398 \\
-4028 \\
20144 \\
0 \\
0 \\
-534 \\
0 \\
-5698 \\
-31805
\end{array}\right] N
$$

Recalling expressions (7.a) and (7.b), $\boldsymbol{u}^{h}$ is interpolated using global shape functions and nodal solutions.

$$
\begin{aligned}
u^{h} & =N_{i} u_{i} \\
v^{h} & =N_{i} v_{i}
\end{aligned}
$$

For $i$ ranging from 1 to 6 .

